

## Physics 727: Assignment 1

(to be submitted by Tuesday, February 13, 2024)

1. Given a square matrix  $A$  of dimension  $n \times n$  with all nonnegative entries, prove that  $A$  is irreducible if  $(I + A)^{n-1}$  has all positive entries.
2. Consider the permutation matrix

$$T = \begin{pmatrix} 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 \end{pmatrix}.$$

- (a) Draw a directed graph showing how the states  $\{|1\rangle, |2\rangle, \dots, |6\rangle\}$  are reconfigured in each time step. Note that there is no branching. Argue that the effect of  $T$  is simply to permute the entries of whatever 6-vector it acts on. You should be able to convince yourself that  $T \simeq (1\ 5\ 3)(2\ 4)(6)$ , which is the permutation expressed in cycle notation.
- (b) Determine whether  $T$  is irreducible. Contrast this with the result for the matrix representation of the permutation  $(1\ 2\ 3\ 4\ 5\ 6)$ .
- (c) Determine whether  $T$  describes a stochastic process that is periodic; if so, determine the period by finding the smallest integer  $k > 1$  such that  $T^k = I$ .
- (d) Are there subspaces of smaller periodicity than what you reported in question 1(d)? If so, what are they?
- (e) Imagine that all the zero entries of  $T$  are replaced by a small positive value  $0 < \epsilon \leq 1/5$  and that all the unit entries are reduced accordingly to maintain the stochastic property:

$$T \rightarrow T_\epsilon = \begin{pmatrix} \epsilon & \epsilon & \epsilon & \epsilon & 1 - 5\epsilon & \epsilon \\ \epsilon & \epsilon & \epsilon & 1 - 5\epsilon & \epsilon & \epsilon \\ 1 - 5\epsilon & \epsilon & \epsilon & \epsilon & \epsilon & \epsilon \\ \epsilon & 1 - 5\epsilon & \epsilon & \epsilon & \epsilon & \epsilon \\ \epsilon & \epsilon & 1 - 5\epsilon & \epsilon & \epsilon & \epsilon \\ \epsilon & \epsilon & \epsilon & \epsilon & \epsilon & 1 - 5\epsilon \end{pmatrix}.$$

Show that  $T_\epsilon$  is aperiodic and irreducible (and hence ergodic).

- (f) Report the long-time equilibrium distribution  $\pi^* = \lim_{k \rightarrow \infty} (T_\epsilon)^k \pi^{(0)}$ .
  - (g) Speculate as to why even an infinitesimal value of  $\epsilon$  is enough to restore ergodicity to the stochastic matrix.
  - (h) What can you say about the timescale for the system to equilibrate and how it depends on  $\epsilon$ ? For instance, if you have two such stochastic systems, one with  $\epsilon = 0.01$  and one with  $\epsilon = 0.02$ , how would the equilibration times compare?
3. Consider the matrix

$$\begin{pmatrix} 1/2 & 1/3 & 1/6 \\ 1/3 & 1/6 & 1/2 \\ 1/6 & 1/2 & 1/3 \end{pmatrix}$$

- (a) List the properties that make  $T$  a doubly stochastic matrix.
- (b) Draw a directed graph showing how the states  $\{|1\rangle, |2\rangle, |3\rangle\}$  are reconfigured in each time step. Label the branching probabilities.
- (c) Determine whether  $T$  is irreducible.

- (d) Determine whether  $T$  describes a stochastic process that is periodic; if so, determine the period by finding the smallest interger  $k > 1$  such that  $T^k = I$ .
- (e) Report the long-time equilibrium distribution  $\pi^* = \lim_{k \rightarrow \infty} T^k \pi^{(0)}$ .
- (f) Show that it takes about 30 time steps for the system to equilibrate from an initial state  $\pi_i^{(0)} = \delta_{i,2}$  to agreement with the final distribution  $\pi^*$  at the 1 in  $10^{16}$  level (double-precision floating point accuracy on a computer).
4. Consider an  $n$ -state system in equilibrium. Each state  $i = 1, 2, \dots, n$  is distributed according to probability  $\pi_i$  and has associated with it an energy  $\varepsilon_i$  and a particle occupancy  $N_i$
- (a) Suppose that the system is isolated from its environment so that the total energy  $U = \sum_i \varepsilon_i \pi_i$  and total particle number  $N = \sum_i N_i \pi_i$  are both held fixed. Augment the entropy  $S = -\sum_i \pi_i \ln \pi_i$  with Lagrange multiplier terms in order to enforce the conservation of energy and particle number. Show that maximizing the entropy results in  $\pi_i = Z^{-1} \exp(-\beta \varepsilon_i + \alpha N_i)$ , where the normalization is given by the partition function  $Z = \sum_i \exp(-\beta \varepsilon_i + \alpha N_i)$ . (More commonly, one sees the expressions  $Z = \sum_i \exp[-\beta(\varepsilon_i - \mu N_i)]$ , with the correspondence  $\alpha = \beta \mu$ .)
- (b) Using the result of part (a), prove that the entropy can be expressed as  $S = \ln Z + \beta U - \alpha N$  (or  $S = \ln Z + \beta U - \beta \mu N$ ).
- (c) Let's treat small variations  $\pi_i \rightarrow \pi_i + \delta \pi_i$ . Show that those variations obey  $\sum_i \delta \pi_i = 0$  and lead to changes  $\delta U = \sum_i \varepsilon_i \delta \pi_i$  and  $\delta N = \sum_i N_i \delta \pi_i$  in the total energy and particle number.
- (d) Demonstrate that such variations in the entropy are related to independent variations in the energy and particle number according to  $\delta S = \beta \delta U - \alpha \delta N$ . Hence, the entropy  $S = S(U, N)$  has partial derivatives  $\beta = \partial S / \partial U$  and  $\alpha = -\partial S / \partial N$  (or  $\mu = -\beta^{-1} \partial S / \partial N$ )
5. Start from the partition function  $Z(\beta, \alpha) = \sum_i \exp(-\beta \varepsilon_i - \alpha N_i)$  in question 4 and define an expectation value operation

$$\langle \cdot \rangle = \frac{1}{Z} \sum_i (\cdot) \exp(-\beta \varepsilon_i - \alpha N_i).$$

For a statistical ensemble in which  $\beta$  and  $\alpha$  are the control parameters, the internal energy and total particle number fluctuate about their average values. Show that such fluctuations are given by

$$\langle (\Delta E)^2 \rangle = \langle E^2 \rangle - \langle E \rangle^2 = \frac{\partial^2 \ln Z}{\partial \beta^2}$$

and

$$\langle (\Delta N)^2 \rangle = \langle N^2 \rangle - \langle N \rangle^2 = \frac{\partial^2 \ln Z}{\partial \alpha^2}.$$

*N.B.:*  $z = e^\alpha$  is often referred to as the fugacity; both  $\alpha$  and  $z$  are dimensionless. The chemical potential  $\mu = \alpha / \beta$  is a rescaled quantity having units of energy (the energy required to add or remove one particle from the equilibrated system). It is more common to express the Boltzmann factor as  $\exp(-\beta \varepsilon_i + \alpha N_i) = \exp[-\beta(\varepsilon_i - \mu N_i)]$ .