GW propagation as a probe for cosmology and beyond

Atsushi Nishizawa (KMI, Nagoya Univ.)



Kobayashi-Maskawa Institute for the Origin of Particles and the Universe

Feb 27 - Mar 2, 2017 StronG BaD workshop @ U of Mississippi

Contents

- 1. Introduction
- 2. GW as a probe for cosmic expansion
 - Hubble const measurement by identifying a host galaxy
- 3. GW as a probe for matter inhomogeneities in the universe
 - angular clustering
 - gravitational lensing

Gravitational Waves

GWs have been detected! (GW150914, GW151226)
 [LIGO Scientific Collaboration 2016]



- BBH merger rate $9 240 \, \mathrm{Gpc}^{-3} \, \mathrm{yr}^{-1}$
- aLIGO is expected to detect more events $\sim 100\,{
 m yr}^{-1}$ This opens new windows to astronomy, cosmology, and gravity.

GW and late-time cosmology



- Cosmic expansion changes distance to a GW source and adds extra GW phase.
- Large scale structure deforms a GW waveform, mainly in amp. (gravitational lensing)
- GW source location clusters and may trace matter dist.

GW as a probe for cosmic expansion

Standard siren (1)

GW from a compact binary can be a cosmological tool to measure distance to a source. [Schutz 1986, Holz & Hughes 2005]

GW phase From observational data, from $L_{\rm gw} = -\frac{dE_{\rm orbit}}{dt}$ $h, f, f \dots$ $\dot{f}(t) \propto \{(1+z)M_c\}^{5/3} f^{11/3}$ $M_z \equiv (1+z)M_c$ **GW** amplitude $h(t) \propto \frac{\{(1+z)M_c\}^{5/3}f^{2/3}}{D_L}$ luminosity distance

Standard siren (2)

In principle, redshift and chirp mass are degenerated.

Assuming the redshift is determined by EM observation of a host galaxy or an EM counterpart,

 M_c is separately determined.

 $z - D_L$ relation is obtained.

Advantages of cosmological-expansion measurement by GWs

- No need of distance ladder.
- Consistency test of la-type SNe observation.
- Accurate cosmological probe (less systematic err)

identifying a host galaxy

[AN 2016]

If a host galaxy for a GW event is uniquely identified, the redshift is obtained from EM obs of the galaxy.

No need to prepare a complete galaxy catalog.

Applicable not only to NS binaries but also to BH binaries.

Ability to identify a host galaxy strongly depends on sky localization error.

Statistical study of parameter estimation

Parameter error estimation with a Fisher information matrix

- log-flat mass distribution (1/m) with $5\,M_\odot < m_1, m_2 < 100\,M_\odot$ and $M < 100\,M_\odot$
- constant merger rate
- 10⁴ nonspinning binaries with S/N > 8
- orbital inclination: uniformly random
- sky position: uniformly random
- phenomenological IMR GW waveform [Khan et al. 2016]
- detector network: aLIGOx2 + aVIRGO (HLV) aLIGOx2 + aVIRGO + KAGRA (HLVK)

Error volume

For each GW event, 3D sky error volume $(\Delta d_L, \Delta \Omega_S)$ is obtained.

of host galaxies in the error volume

$$N_{\text{host}} \equiv n_{\text{gal}} \left\{ V(d_{L,\text{max}}) - V(d_{L,\text{min}}) \right\} \frac{\Delta \Omega}{4\tau}$$

$$n_{\rm gal} = 0.01 \, \rm Mpc^{-3}$$

 $d_{L,\min}$ $d_{L,\max}$ galaxy number density (covering 90% total luminosity in B-band)

source selection	source catalog	$70{ m Gpc}^{-3}{ m yr}^{-1}$	$30{ m Gpc}^{-3}{ m yr}^{-1}$	$10{ m Gpc}^{-3}{ m yr}^{-1}$
HLV (SNR>8, $z < 1$)	10000	$4512~{\rm yr}^{-1}$	$1934 { m yr}^{-1}$	$645~{ m yr}^{-1}$
HLV (SNR>8, $z < 1$), $N_{\text{host}} < 1$	49	$22~{ m yr}^{-1}$	$9~{ m yr}^{-1}$	$3~{ m yr}^{-1}$
HLVK (SNR>8, $z < 1$)	10000	$5122~{ m yr}^{-1}$	$2195~{ m yr}^{-1}$	$732~{ m yr}^{-1}$
HLVK (SNR>8, $z < 1$), $N_{\text{host}} < 1$	59	$30~{ m yr}^{-1}$	$13~{ m yr}^{-1}$	$4~{ m yr}^{-1}$

 $\Delta\Omega_{
m S}$



measurement precision of Hubble const



Can be reduced to 1%. [Tuyenbayev et al. 2017]

H₀ value discrepancy problem

[Planck Collaboration 2013]

There is still some discrepancy between observations of H_0 .

GW from BH binaries allows us to measure H_0 at precision of 0.8%.

Importantly, GW obs is completely independent way to measure cosmic expansion.



H₀ value discrepancy problem

[Planck Collaboration 2013]

There is still some discrepancy between observations of H_0 .

GW from BH binaries allows us to measure H_0 at precision of 0.8%.

Importantly, GW obs is completely independent way to measure cosmic expansion.



GW as a probe for matter inhomogeneities in the universe

Beyond background cosmology

- With a standard siren, one can measure luminosity distance to a source directly. Combining the luminosity distance with redshift information, cosmic expansion is measured.
- models for cosmic accelerating expansion dark energy (scalar field etc.) vs modification of gravity

The problem is that most models can mimic the LCDM model as a special case by tuning their model parameters.

To discriminate models, need to go to a perturbative level. (matter clustering, gravitational lensing, etc.)

angular clustering

cross-correlating sky maps





galaxy survey (2MASS redshift survey)

- GW events
- How strongly are GW events correlated with galaxies?
- What kind of galaxies are associated with BBH?

Х

angular cross-power spectrum [Namikawa, AN, Taruya 2016]

cross-correlation between the probes, GW (s) and galaxy (g)

$$C_{\ell}^{sg} = 4\pi \int_{0}^{\infty} d\ln k \int_{0}^{\infty} d\chi \, j_{\ell}(k\chi) \int_{0}^{\infty} d\chi' \, j_{\ell}(k\chi') \times W^{s}(k,\chi) W^{g}(k,\chi') \Delta_{m}(k;\chi,\chi')$$

weight function

matter density power spectrum

$$W^{s}(\chi) = \frac{dn_{\rm BH}}{d\chi}(\chi) \underbrace{b_{\rm BH}(\chi)}_{\downarrow} \qquad \text{for GW}$$

$$W^{g}(\chi) = \frac{dn_{\rm gal}}{d\chi}(\chi) \underbrace{b_{\rm gal}(\chi)}_{\downarrow} \qquad \text{for galaxy}$$

$$(\text{ustering strength of galaxy})$$

power spectrum: GW x galaxy

aLIGO x2 + aVIRGO observations (at design sensitivity) & Pan-STARRS



detection significance of clustering

$$\alpha_{sg} = \alpha_{sg}^0 \left(\frac{b_{\rm BH,0}}{1.5}\right) \left(\frac{T_{\rm obs} \dot{n}_0}{3 \times 100 \,\rm Gpc^{-3}}\right)^{1/2}$$

 $\begin{array}{ll} \mbox{GW x Euclid} & \mbox{GW x Pan-STARRS} \\ \alpha^0_{sg} = 3.6 & \alpha^0_{sg} = 4.5 \end{array}$

BBH clustering $b_{BH,0}$ can be detected unless BBH merger rate is so small.

If $b_{\rm BH,0} \approx b_{\rm gal,0}$, BBH are likely to trace a baryon distribution and star formation.

If not, a nonstandard scenario (e.g. PBH) may be preferred.

gravitational lensing

Gravitational lensing of GW [Wang, Stebbins & Turner 1996, Holz & Wald 1998]



- GW traces its null geodesic and is lensed by galaxies and galaxy clusters.
- Source is a compact binary.
 No shear (too small image), but "brightness" of GW is magnified or demagnified.

- Apparent luminosity distance $D(\mathbf{x}) = \bar{D} \left\{ 1 + \kappa(\mathbf{x}) \right\}$

"magnification"

anisotropy of luminosity distance

[Namikawa, AN, Taruya 2016]

deviation of luminosity distance from the averaged one in i-th distance bin



$$\begin{split} \hat{s}_i(\Omega) &= \frac{d_i(\Omega) - d_i}{\hat{d}_i} & \text{average number} & \text{i-th distance} \\ &= \frac{1}{\bar{d}_i} \int_{D_i^{\min}}^{D_i^{\max}} dD \ \bar{n}(D) \{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\min}}^{D_i^{\max}} dD \ \bar{n}(D) \{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\min}}^{D_i^{\max}} dD \ \bar{n}(D) \{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\min}}^{D_i^{\max}} dD \ \bar{n}(D) \{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\min}}^{D_i^{\max}} dD \ \bar{n}(D) \{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\min}}^{D_i^{\max}} dD \ \bar{n}(D) \{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\min}}^{D_i^{\max}} dD \ \bar{n}(D) \{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\min}}^{D_i^{\max}} dD \ \bar{n}(D) \{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\min}}^{D_i^{\max}} dD \ \bar{n}(D) \{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\min}}^{D_i^{\max}} dD \ \bar{n}(D) \{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\max}}^{D_i^{\max}} dD \ \bar{n}(D) \{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\max}}^{D_i^{\max}} dD \ \bar{n}(D) \{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\max}}^{D_i^{\max}} dD \ \bar{n}(D) \{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\max}}^{D_i^{\max}} dD \ \bar{n}(D) \{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\max}}^{D_i^{\max}} dD \ \bar{n}(D) \{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\max}}^{D_i^{\max}} dD \ \bar{n}(D) \{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\max}}^{D_i^{\max}} dD \ \bar{n}(D) \{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\max}}^{D_i^{\max}} dD \ \bar{n}(D) \{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\max}}^{D_i^{\max}} dD \ \bar{n}(D) \{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\max}}^{D_i^{\max}} dD \ \bar{n}(D) \{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\max}}^{D_i^{\max}} dD \ \bar{n}(D) \{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\max}}^{D_i^{\max}} dD \ \bar{n}(D) \left\{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\max}}^{D_i^{\max}} dD \ \bar{n}(D) \left\{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\max}}^{D_i^{\max}} dD \ \bar{n}(D) \left\{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\max}}^{D_i^{\max}} dD \ \bar{n}(D) \left\{ \delta(\mathbf{x}) + \gamma(D) \kappa(\mathbf{x}) \} \\ &= \hat{d}_i \int_{D_i^{\max}}^{D_i^{\max}} dD \ \bar{n}(D) \left\{ \delta(\mathbf{x}) + \gamma(D) \kappa(D) \kappa(D) \right\} \\ &= \hat{d}_i \int_{D_i^{\max}}^{D_i$$



cosmological implications

- non-Gaussianity of large-scale structure with ET $\sigma(f_{\rm NL}) \approx 0.54 \qquad \begin{array}{c} {\rm comparable \ or \ better} \\ {\rm than \ Euclid} \end{array}$
- cross-correlation of clustering

GW (ET) x PlanckSNR ~ 31GW (ET) x CMB stage IVSNR ~ 43

- cross-correlation of weak lensing
 GW (ET) x Euclid SNR ~ 16
- A lot of applications of GW observations to cosmology

Open questions

 With BBH observed with aLIGOx2 + aVIRGO, Hubble constant can be measured at 1% level.

How realistic is this method? Any more systematic error?

• Angular clustering of GW from BBH gives information about what type of galaxies are associated with them.

Any robust prediction in astrophysical side?

• Gravitational lensing of GW offers many cosmological applications by correlating with CMB and galaxy surveys.

Sensitivity to cosmological parameters? Any systematic bias?