

**The Liouville Theorem**

- *Statement:* The value of the distribution function along a trajectory in phase space is a constant of the motion, or  $\rho$  evolves like an incompressible fluid. Technically, the statement that

$$\frac{d\rho}{dt} = \frac{\partial\rho}{\partial t} + \sum_{\text{dof}} \left( \frac{\partial\rho}{\partial q^i} \dot{q}^i + \frac{\partial\rho}{\partial p_i} \dot{p}_i \right) = 0 .$$

- *Proof:* The time rate of change of the probability of a region  $R$  in  $\Gamma$  must equal the flux of probability through its surface, which means that  $\rho$  must satisfy an equation of continuity,

$$0 = \frac{\partial\rho}{\partial t} + \vec{\nabla} \cdot (\vec{v}\rho) = \frac{\partial\rho}{\partial t} + \left( \frac{\partial\dot{q}^i}{\partial q^i} + \frac{\partial\dot{p}_i}{\partial p_i} \right) \rho + \left( \frac{\partial\rho}{\partial q^i} \dot{q}^i + \frac{\partial\rho}{\partial p_i} \dot{p}_i \right) ,$$

where  $\vec{v} = (\dot{q}^1, \dot{q}^2, \dots, \dot{q}^{3N}; \dot{p}_1, \dot{p}_2, \dots, \dot{p}_{3N})$ , and the terms in the first set of parentheses cancel by virtue of the Hamilton equations of motion.

**General Considerations on Equilibrium Distribution Functions**

- *Recall:* The distribution function  $\rho(q, p; t)$  is a constant of the motion.
- *Goal:* Find a general form for  $\rho(q, p)$  for an equilibrium system.
- *Possible situations:* Two situations in which a system can reach equilibrium are if it is isolated (in which case we will assume the value of the energy is known), and if it is in thermal equilibrium with a bath at fixed temperature, with which it can exchange energy and possibly particles.

**The Microcanonical Distribution**

- *Idea:* The distribution function describing an isolated system with fixed value for the energy.
- *Distribution function:* The energy of the system is conserved, so if it known to have a certain value  $E$ ,

$$\rho(q, p) = \text{constant} \times \delta(H(q, p) - E) , \quad \text{the microcanonical distribution ,}$$

which incorporates the principle of equal a priori probabilities (otherwise the “constant” could have been any function of all  $6N - 1$  constants of the motion. [To avoid the use of this axiom, we could also have used the large-system approximation as in the thermal equilibrium case, but that would have had an intrinsic contradiction, since the log of this  $\rho$  is not linear in  $H$ .])

**Relevant Sections:** Phys 731; Chandler; Halley, second half of Ch 1; Reif; Schwabl.

**Related Topics**

- *Additive constants of the motion:* @ J M Amigo & H Reeh, *Annalen der Physik* **36** (1988) 929-937.
- *Non-extensive statistics:* What changes with respect to the treatment here?
- *Other possible distribution functions:* Which ones have been used?